

Random Systems

- Real systems are usually "dirty": many impurities, etc.
- Model this using random couplings (= disorder)

Random Ising model: $\mathcal{H} = - \sum_{\langle ij \rangle} J_{ij} S_i S_j$
random variables drawn from, say, a Gaussian distribution.

- Disordered systems have a very rich phenomenology: spin glasses, etc. Here, we focus on a simple question:

Does disorder/randomness modify the critical behavior (universality class) of a transition?

↳ is weak disorder relevant at a given RG fixed point?

A Replica Trick

Important to distinguish:

* "Annealed" disorder: impurities are dynamical variables like the spins in our Ising model, and should be averaged over in the partition function: $Z_{\text{annealed}} = \overline{Z}$ (average over J_{ij})

* "Quenched" Disorder: impurities are "frozen", Z depends on J_{ij}
→ loose translation invariance.

However, for a thermodynamic system: $\frac{F}{N} \approx \overline{\frac{F}{N}}$ (we say that the free energy is self averaging)

This is because $F = \sum_{\text{subsystems } i} F_i$ in the thermodynamic limit.
 $\frac{F}{N} = \left(\sum_{\text{subsystems } i} \frac{F_i}{N_{\text{subsystem}}} \right) / \# \text{ subsystems}$

Consider quenched disorder: we need to average F : compute \bar{F}
 → as usual, all physical quantities follow from taking derivatives.

Problem: computing $\bar{F} = -\overline{\log Z}$ is hard!

Replica trick: use formula: $\log Z = \lim_{n \rightarrow 0} \frac{Z^n - 1}{n}$

⇒ Z^n = partition function of n copies ("replicas") of the system for $n \in \mathbb{N}$
 and \bar{Z}^n usually much easier to compute! (for n integer)
 ⇒ then analytically continue to $n \rightarrow 0$ (can be subtle! may not commute with thermodynamic limit!)

(B) Harris Criterion

Consider a clean model: $\beta \mathcal{H} = \beta \mathcal{H}^* + t \int d^d x \phi_T(x)$
 ↑ RG fixed point ↑ thermal perturbation
 $\Delta_t = d - \nu^{-1}$

Coarse graining: impurities couple to energy density $\epsilon(x) \sim \phi_T(x) + \dots$

⇒ $\beta \mathcal{H} = \beta \mathcal{H}^* + \int d^d x t(x) \phi_T(x) + \dots$

Replica trick: $\bar{Z}^n = \int D\vec{\phi} e^{-\sum_{a=1}^n \beta \mathcal{H}[\phi_a]}$
 disorder = random variables not correlated in space:
 $\delta t(x) \delta t(x') = \sigma_T^2 \delta(x-x')$
 with $\delta t(x) = t(x) - \bar{t}$
 degrees of freedom in the model = ϕ
 $\vec{\phi} = (\phi_1, \dots, \phi_n)$ in replicated theory

⇒ $\bar{Z}^n = \int D\vec{\phi} e^{-\beta \sum_{a=1}^n \mathcal{H}^*[\phi_a]} e^{-\bar{t} \int d^d x \sum_a \phi_a + \frac{\sigma_T^2}{2} \int d^d x \sum_a \phi_a^2 + \dots}$
 cumulant expansion disorder strength used $\delta t(x) \delta t(x') = \sigma_T^2 \delta(x-x')$

* 1st term = trivial: F just shifts T_c

* 2nd term: $\sigma_F^2 = \overline{F^2} - (\overline{F})^2$: Variance of the disorder (= 0 if clean system)

\Rightarrow effective interaction between replicas: $-\frac{\sigma_F^2}{2} \int d\vec{r} \sum_{a,b=1}^n \phi_F^a(\vec{r}) \phi_F^b(\vec{r})$

\Rightarrow $a=b$ terms: same replica, $(\phi_F^a)^2$ not a scaling operator in general
but $(\phi_F^a)^2 \simeq c \phi_F^a + \dots \Rightarrow$ effective shift in T_c , doesn't change critical exponents.

\Rightarrow $a \neq b$ terms: $\mathcal{O}(\vec{r}) = \sum_{a \neq b} \phi_F^a(\vec{r}) \phi_F^b(\vec{r})$ relevant? Scaling dimension?
 $S_0 = \beta \sum_{a=1}^n \mathcal{H}^*[\phi_a]$ decoupled replicas

$\Rightarrow \langle \mathcal{O}(\vec{r}) \mathcal{O}(\vec{r}') \rangle = 2n(n-1) \langle \phi_F^a(\vec{r}) \phi_F^a(\vec{r}') \rangle^2 \sim \frac{2n(n-1)}{|\vec{r} - \vec{r}'|^{4\Delta_F}}$

this yields: $\Delta_{\mathcal{O}} = 2\Delta_F$
 $\gamma_{\mathcal{O}} = d - 2\Delta_F = \frac{2}{\nu} - d$

(Weak) disorder is irrelevant if $\gamma_{\mathcal{O}} < 0 \Leftrightarrow d\nu > 2$ Harris Criterion

correlation length exponent of the clean fixed point

• if this inequality is satisfied: disorder irrelevant, and just shifts T_c (but does not change critical exponents).

• if $d\nu < 2$: new fixed point / physics in the presence of disorder