

Standard Model 4

Note Title

11/19/2009

Pion/Kaons in QCD

$$m_q \rightarrow 0$$

$$\psi_L \rightarrow L \psi_L$$

$$\psi_R \rightarrow R \psi_R$$

don't see partners

- Symmetry hidden \Rightarrow Goldstone bosons

$$P + P + \pi(q=0)$$

Dynamical Symmetry Breaking

$$\Rightarrow \mathcal{L} = \frac{F^2}{4} \text{Tr}(\partial_\mu U \partial^\mu U^\dagger) + \dots$$

interesting

External source technique

$$\psi = \begin{pmatrix} u \\ d \\ s \end{pmatrix}$$

$$\mathcal{L}_{QCD} = -\frac{1}{4} F^2 + \bar{\psi} (i \not{D} - m) \psi \quad \leftarrow \text{SU(3) in color}$$

$$\hat{D} = \left(\not{\partial} + ig \frac{\lambda^a G^a}{2} \right) \mathbb{1} \quad \leftarrow \text{in flavor}$$

$$\bar{\psi}_m \psi = \bar{\psi}_L m \psi_L + \bar{\psi}_R m \psi_L \Rightarrow m = \begin{pmatrix} m_u & & \\ & m_d & \\ & & m_s \end{pmatrix}$$

Add external sources

$$L_m, R_m, S, P \leftarrow 3 \times 3 \text{ in flavor space}$$

$$S = S^0 + t^A S^A$$

$$\leftarrow \text{SU(3) matrix in flavor SU(3)}$$

$$\mathcal{L} = -\frac{1}{4} F^2 + \bar{\Psi}_L i \not{D}_L \Psi_L + \bar{\Psi}_R i \not{D}_R \Psi_R + \bar{\Psi}_L (S + ip) \Psi_R + \bar{\Psi}_R (S - ip) \Psi_L$$

$$\not{D}_L = \left(\partial_\mu + ig \frac{\sigma^a}{2} G^a \right) \not{\mathbb{1}} + i \not{\ell}_\mu(x)$$

$$\not{D}_R = \left(\partial_\mu + ig \frac{\sigma^a}{2} G^a \right) \not{\mathbb{1}} + i \not{r}_\mu(x)$$

Pure QCD $\Rightarrow \ell_\mu = r_\mu = 0$, $S = M$, $p = 0$

QED + photon $\ell_\mu = r_\mu = e Q A_\mu$, $S = m$, $p = 0$
 $\tau \begin{pmatrix} 2/3 & -1/3 & -1/3 \end{pmatrix}$

QCD + W $r_\mu = 0$, $\ell_\mu = \frac{g_2}{\sqrt{2}} \left(W_\mu^+ \begin{pmatrix} 0 & V_{ud} & V_{us} \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} + h.c. \right)$

Currents

$$\frac{\partial \mathcal{L}}{\partial S^0} = \bar{\psi} \psi$$

$$\frac{\partial \mathcal{L}}{\partial l_\mu^a} = \bar{\psi}_L \gamma_\mu \tau^a \psi_L$$

$$\frac{\partial \mathcal{L}}{\partial A^\mu} = \bar{\psi} \gamma^\mu Q \psi = J_\mu^{em}$$

Enhanced Symmetry

$$\psi_L \rightarrow L \psi_L$$

$$\psi_R \rightarrow R \psi_R$$

$$+ (S+ip) \rightarrow L (S+ip) R^\dagger$$

$$\bar{\psi}_L (S+ip) \psi_R \text{ - invariant}$$

exactly invariant

make it $SU(3)_L \times SU(3)_R$ gauge symmetry $\mathcal{L}_m \rightarrow L (\mathcal{L}_m + L^\dagger \partial_\mu L) L^\dagger$

$$\bar{\psi}_L \not{\partial} \psi_L \Rightarrow \bar{\psi} L^\dagger (\not{\partial} + i L (\mathcal{L}_m + L^\dagger \not{\partial} L) L^\dagger) L \psi_L$$

$$= \bar{\psi} \left[\cancel{(L^\dagger \not{\partial} L)} + i \not{\partial} + i \not{\partial} \cancel{- i L^\dagger \not{\partial} L} \right] \psi_L = \bar{\psi}_L \not{\partial} \psi_L$$

$$D_{mL} \rightarrow L D_m L^\dagger$$

$$\psi_L \rightarrow L \psi_L, \quad h_m \rightarrow L [\quad] L^+$$

$$r_n \rightarrow R [\quad] R^+$$

$$(s+ip) \rightarrow L (s+ip) R^+$$

$$[D_{mL}, D_{nL}] = \partial_m h_n - \partial_n h_m + [h_m, h_n] = L_{mn}$$

$$L_{mn} \rightarrow L L_{mn} L^+$$

$$R_{mn} \rightarrow R R_{mn} R^+$$

Go to path integral

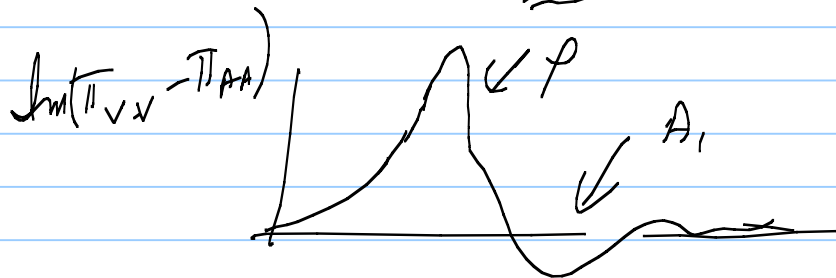
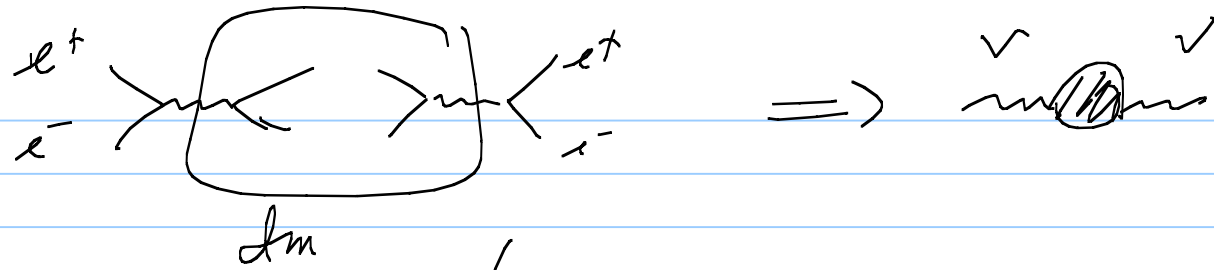
$$Z[\sigma, \rho, l, r] = \int dA_\mu d\psi d\bar{\psi} e^{i \int d^4x \mathcal{L}(\psi, A, \sigma, \rho, l, r)}$$

$$\langle 0 | \bar{\psi} \psi | 0 \rangle = \frac{i \delta Z}{Z \delta \sigma^0}$$

→ signal of chiral symmetry breaking
(eigen of ψ in SSB)

$$\langle 0 | \bar{\psi} \frac{t^a}{2} \gamma_5 (1 + \gamma_5) \psi(x) \quad \bar{\psi} \frac{t^b}{2} \gamma_5 (1 + \gamma_5) \psi(y) | 0 \rangle = \frac{(-i)^2 \delta Z}{Z \delta \sigma^a(x) \delta \sigma^b(y)}$$

$$\frac{\delta}{\delta l \delta r} = \langle \bar{\psi} \gamma^m (1 + \gamma_5) \psi \quad \bar{\psi} \gamma_n (1 - \gamma_5) \psi(y) \rangle = 0 \quad \text{if no SSB}$$



Effective \mathcal{L}

- only d.o.f π, κ

Match

$$\begin{aligned} Z[\sigma, l, r] &= \int dA d\psi d\bar{\psi} \mathcal{L} \\ &= \int d\pi e^{i \int d^4x \mathcal{L}_{\text{eff}}(\pi, \sigma, p, l, r)} \end{aligned}$$

$i \int d^4x \mathcal{L}_{\text{eff}}$
 \uparrow exact $SU(3)$ gauge sym

Global $SU(3)$

$$U = \exp i \frac{t^a \phi^a}{F}$$

$$U \rightarrow LUR^\dagger$$

$$\text{Tr}(\partial_\mu U \partial^\mu U^\dagger) + -$$

Gauge symm.

$$D_\mu U \rightarrow L (D_\mu U) R^\dagger$$

Solutions

$$D_\mu U = \partial_\mu U + i l_\mu U - i U R_\mu$$

$$L_{\mu\nu}, R^{\mu\nu}$$

$$\mathcal{S} + i\rho \Rightarrow \chi = 2B_0 (\mathcal{S} + i\rho) \Rightarrow \chi \rightarrow L \chi R^\dagger$$

$$\partial_\mu \sim E$$

$$l_\mu \sim E$$

$$L_{\mu\nu} \sim E^2$$

and $\chi \sim E^2$

after $S \rightarrow m$

$$m_{\pi}^2 = 2B_0 m$$

↑ →

Constructing \mathcal{L}_{eff}

$$E^0, \quad U^\dagger U = 1$$

$$E^2, \quad \text{Tr}(D_\mu U (D_\mu U)^\dagger), \quad \text{Tr}(X^\dagger U + U^\dagger X)$$

$$\mathcal{L}_2 = \frac{F^3}{4} \text{Tr}(D_\mu U D^\mu U^\dagger) + \frac{F^2}{4} \text{Tr}(X^\dagger U + U^\dagger X) \quad (F, B_\mu)$$

↑ 2 B.S

$\mathcal{O}(E^4)$

$$\mathcal{L} = L_1 \left[\text{Tr} D_\mu U D^\mu U^\dagger \right]^2 \dots \text{Tr}(X U + U^\dagger X) \text{Tr}(D_\mu U D^\mu U^\dagger)$$

+ $L_{10} \text{Tr}(L_{\mu\nu} U R^{\mu\nu} U^\dagger)$

↑ + 10 terms

Explore 2

$$1) \langle 0 | \bar{\psi} \psi | 0 \rangle = \frac{\delta Z}{\delta S^0} = \overset{\text{Tree level (no loops)}}{-F^2 B_0} \quad B_0 = -\frac{1}{F} \langle 0 | \bar{\psi} \psi | 0 \rangle$$

2) π & K masses

$$M_\pi^2 = B_0 (m_u + m_d) = 2 B_0 \bar{m}$$

$$M_{K^+}^2 = B_0 (m_u + m_s)$$

$$M_{K^0}^2 = B_0 (m_d + m_s)$$

$$M_{\eta_8}^2 = B_0 \frac{1}{3} (m_u + m_d + 4m_s) = \frac{4}{3} M_{K^+}^2 - \frac{1}{3} M_\pi^2$$

$$\frac{M}{M_s} = \frac{M_{II}^2}{2M_K^2 - M_{II}^2} = \frac{1}{2.6}$$

$$\frac{M_d - M_h}{M_d + M_h} = \frac{M_{K^0}^2 - M_{K^+}^2}{M_{II}^2} = 0.29 \quad \Rightarrow \quad \frac{M_h}{M_d} = 0.55$$

↙ assume B+M

Complete U(1)_A story

IF we had U(3)_A DSB \Rightarrow of G.B.

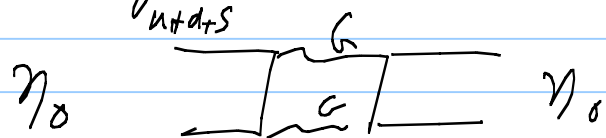
Then masses

$$M^2 = B, \begin{pmatrix} 2\hat{m} & 0 & 0 & 0 \\ 0 & M_S + \hat{m} & 0 & 0 \\ 0 & 0 & \frac{2}{3}(2M_S + \hat{m}) & \frac{2\sqrt{2}}{3}(\hat{m} - M_S) \\ 0 & 0 & \frac{2\sqrt{2}}{3}(\hat{m} - M_S) & \frac{2}{3}(M_S + 2\hat{m}) \end{pmatrix}$$

$$\Rightarrow M_a^2 = M_{\pi}^2, \quad M_b^2 = 2M_K^2 - M_{\pi}^2$$

U(1) problem state degenerate with pions

After $U(1)_A$ anomaly



$$M = \begin{pmatrix} & & & m_s - \tilde{m} \\ & & & + \epsilon \\ & & \text{same} & \\ & & & m_s - \tilde{m} \end{pmatrix}$$

$\epsilon \sim \mathcal{O}(1 \text{ GeV})$

