

# Renormalization 3

Note Title

4/6/2010

## Renorm

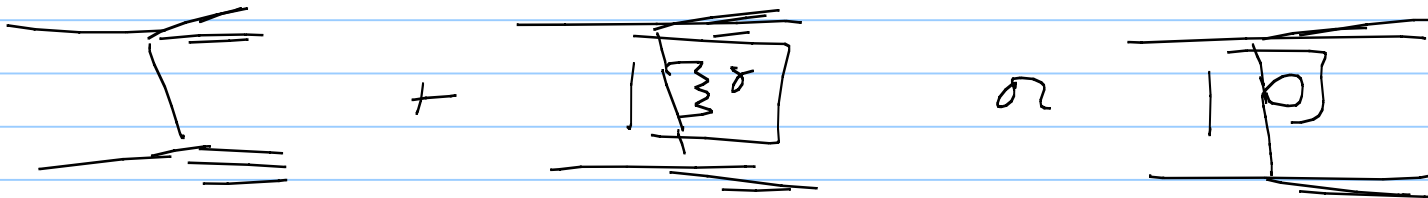
- 1) Coupling constant  $\lambda, e, \dots$
- 2) Mass renorm. } Self energy diagrams
- 3) Wavefunction renorm.

# Propagator

$$\frac{i}{p^2 - m^2 + i\epsilon} = \frac{i}{E^2 - \vec{p}^2 - m^2 + i\epsilon} = i \left[ \frac{P}{E^2 - \vec{p}^2 - m^2} - i\pi \delta(E^2 - \vec{p}^2 - m^2) \right]$$

↙ off shell
↘ on shell

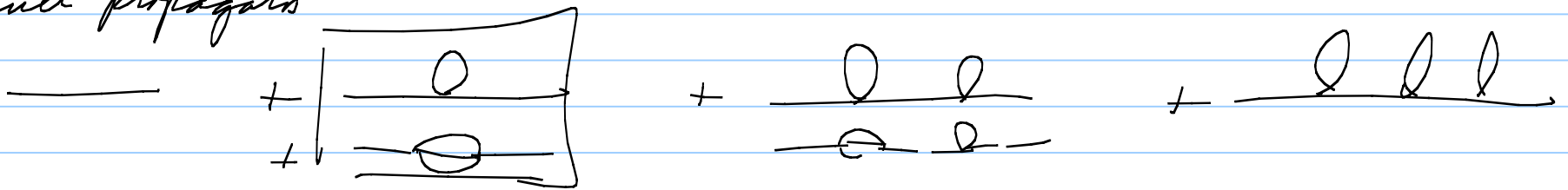
# Interactions



# Self energies

$$-i \Sigma(p) \equiv \text{[Self-energy diagrams]} \quad \text{or} \quad \text{[Self-energy diagrams]} \quad \left( \text{drop external factors } \bar{u}( ) u \right)$$

# Full propagator



$$\begin{aligned}
 & \frac{i}{p^2 - m_0^2 + i\epsilon} + \frac{i}{p^2 - m_0^2 + i\epsilon} - i \Sigma(p) \frac{i}{p^2 - m_0^2 + i\epsilon} + \frac{i}{p^2 - m_0^2 + i\epsilon} - i \Sigma \frac{i}{p^2 - m_0^2 + i\epsilon} - i \Sigma \frac{i}{p^2 - m_0^2 + i\epsilon} + \dots \\
 & = \frac{i}{p^2 - m_0^2 - \Sigma(p) + i\epsilon}
 \end{aligned}$$

Now pole occurs at

$$p^2 - m_0^2 - \Sigma(p) = 0$$

In Pert theory

$$\Sigma(p^2) = \Sigma(m^2) + (p^2 - m^2) \Sigma'(m^2) + \dots$$

$$m^2 = m_0^2 + \Sigma(m^2) = \text{physical mass} \quad \leftarrow \text{pole}$$

Wavefunction renorm

$$\frac{i}{p^2 - m_0^2 - \Sigma(m^2) - \underbrace{(p^2 - m^2) \Sigma'(m^2)}} = \frac{i}{(p^2 - m^2)(1 - \Sigma'(m))} = \frac{iZ}{p^2 - m^2 + i\epsilon}$$

Logic

$$i D_F(x-y) = \langle 0 | T \phi(x) \phi(y) | 0 \rangle$$

$$(\not{D} + m^2) i D_F(x-y) = \langle 0 | [\phi(x), \not{\partial}_0 \phi(y)]_{x_0=y_0} | 0 \rangle \delta(x_0 - y_0) = i \delta^4(x-y)$$

~~\*\*\*~~ defining property

↑ no factor of Z

Need to readjust norm

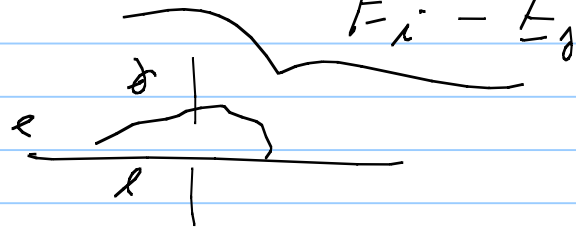
$$\phi_r = Z^{1/2} \phi$$

^  $\phi_r$  will have correct propagator

In QM

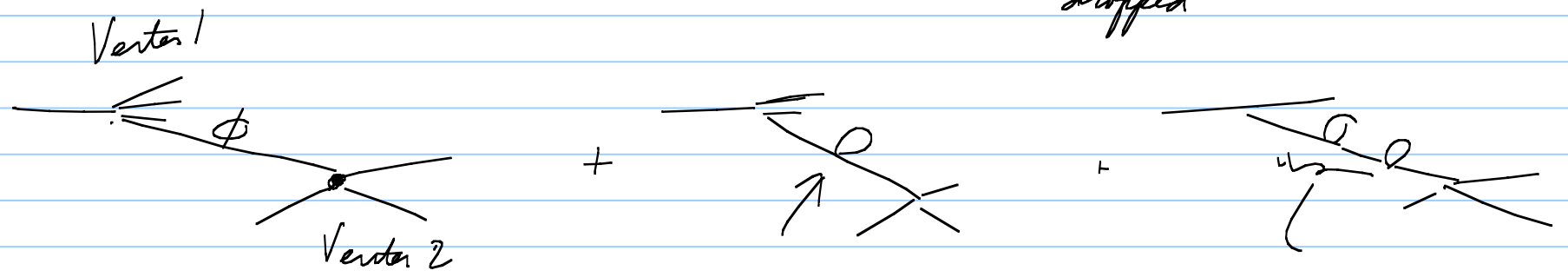
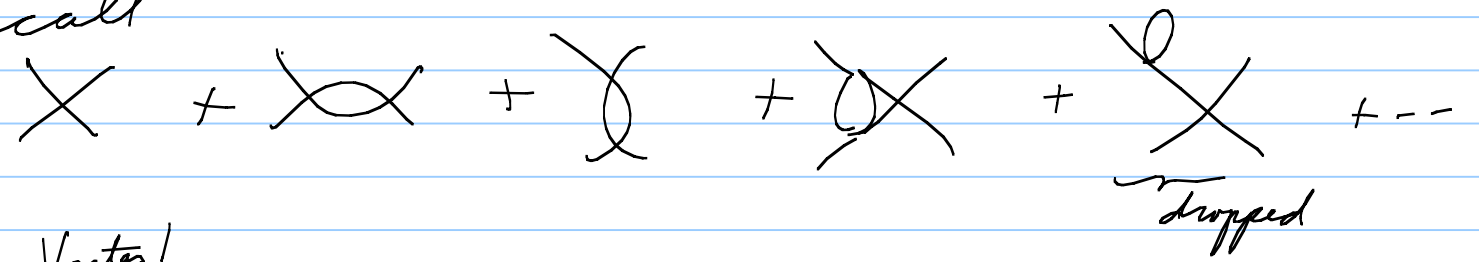
$$|\psi_i\rangle = N \left[ |\psi_{i0}\rangle + \sum_{j \neq i} |\psi_{j0}\rangle \frac{\langle \psi_{j0} | V | \psi_{i0} \rangle}{E_i - E_j} \right]$$

$$\langle \psi_i | \psi_i \rangle = \underline{\hspace{2cm}}$$



# Logic for dropping Self energies

Recall



net effect is just going to physical mass

On shell part  $(\text{Vertex 1}) \sim \pi \delta(p^2 - M_{\text{phys}}^2) \sim (\text{Vertex \#2})$

↑ same effect of  $Z$

Formal technique - mass renorm by itself

- counterterm method

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \phi)^2 - \frac{m_0^2}{2} \phi^2 - \frac{\lambda}{4} \phi^4$$

Express in term of physical mass

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \phi)^2 - \frac{m^2}{2} \phi^2 - \frac{\lambda}{4} \phi^4 + \underbrace{\frac{\delta m^2}{2} \phi^2}_{\text{Treat as perturbations}}$$

$\swarrow m^2 - m_0^2$

- Feynman rule  $\text{---} \overline{\text{---}} \text{---} = i \delta m^2$

Calculate with physical mass

$$-iM = \cancel{X} + \cancel{Y} + \cancel{Z} + \cancel{W} + \cancel{V} + \cancel{U}$$

$$-i\Sigma(p^2) = \text{---} + \text{---} + \text{---}$$

Propagator

$$= \frac{i}{p^2 - M^2 + i\epsilon} + \frac{1}{p^2 - m^2} \begin{matrix} -i\epsilon \\ i\epsilon \end{matrix} \frac{i}{p^2 - m^2} + \frac{i}{p^2 - m^2} i\delta m^2 \frac{1}{p^2 - m^2}$$

$$= \frac{i}{p^2 - m^2 - \underbrace{\Sigma(p) + \delta m^2}_{\text{physical mass}} + i\epsilon}$$

$$\delta m^2 = \Sigma(m^2)$$

⇒ all predictions use physical mass  
- drop external self energies



## General procedure

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \phi)^2 - \frac{M_0^2}{2} \phi^2 - \frac{\lambda_0}{4} \phi^4$$

$$1) \phi_R \equiv Z_\phi^{-1/2} \phi \quad \text{or} \quad \phi = Z_\phi^{1/2} \phi_R$$

$$\mathcal{L} = Z_\phi \frac{1}{2} (\partial_\mu \phi_R)^2 - \frac{M_0^2 Z_\phi}{2} \phi_R^2 - \frac{\lambda_0 Z_\phi^2}{4} \phi_R^4$$

2.) use correct mass + coupling constant  $\Rightarrow$  counterterm

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \phi_R)^2 - \frac{m^2}{2} \phi_R^2 - \frac{\lambda}{4} \phi_R^4 + \mathcal{L}_{ct}$$

$$\mathcal{L}_{ct} = \frac{1}{2} (\delta Z) (\partial_\mu \phi_R)^2 + \frac{\delta m^2}{2} \phi_R^2 - \frac{\delta \lambda}{4} \phi_R^4$$

$$\hat{\delta Z} = Z_\phi - 1$$

3)  $\mathcal{K}_{ct}$  as perturbation

$$\text{---} \otimes \text{---} + i (\delta m^2 + \delta Z(p^2 - m^2))$$

$$\text{---} \otimes \text{---} = -6i \delta \lambda$$

4) Renormalize propagator

$$\text{---} = \text{---} + \text{---} \circ \text{---} + \text{---} \times \text{---} + \dots$$

$$= \frac{i}{p^2 - m^2 - \Sigma(p) + (\delta m^2 + \delta Z(p^2 - m^2))}$$

$$\uparrow \Sigma(m^2) + (p^2 - m^2) \Sigma'(m^2) + \mathcal{O}((p^2 - m^2)^2)$$

right mass + norm if

$$\delta m^2 = \Sigma(m^2)$$

$$\delta Z = (Z_4 - 1) = \Sigma'(m^2) = \left. \frac{\partial \Sigma(p^2)}{\partial p^2} \right|_{p^2 = m^2}$$

5) Coupling constant

$$\delta m^2 = Z_4(m^2 - M_0^2)$$

$$\delta \lambda = \lambda_0 Z_4^2 - \lambda$$

↑  
new

$$-iM \Big|_{\text{renorm}} = \text{tree} + \text{loop} + \dots + \text{loop with } \delta \lambda$$

=  $-6i\lambda$  at the renorm point  $\Rightarrow \delta \lambda$

## 2 Procedures

1) BPH - counterterms

2) Conventional or direct

- work with  $m_0, \lambda_0$  ---

- new rule  $Z_\phi^{1/2}$  for each external state

$$A_n = \lambda_0 Z_\phi^4 \underbrace{\text{diagram}}_{\sqrt{\lambda}}$$

The diagram is a loop with two external lines, crossed out with a large 'X'.

# "Renormalizable" vs "Nonrenormalizable" theories

Any theory can be renormalized (caveats)

- absorb divergence in parameter of  $\mathcal{L}$
- make physical predictions using measured parameters

Caveat: "Whatever is not forbidden is required"

- can't drop terms without reason

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \phi)^2 - \frac{\lambda}{4} \phi^4 \quad (\text{no mass term})$$

$$\underline{0} = \mathcal{E}(p) \Rightarrow \int m^2 \text{ required for renorm}$$

} Not correct

but  $\mathcal{L} = \frac{1}{2} (\partial_\mu \phi)^2 - \frac{\lambda}{4} (\partial_\mu \phi \partial^\mu \phi)^2 + \dots - (\partial_\mu \phi \partial^\mu \phi)^4 + \dots$  is OK to drop mass

Shift Symmetry  $\phi \rightarrow \phi + c$   
 $m^2 \phi^2$  breaks symmetry

"Renormalizable" is separate classification

↑ infinities only in a small # of couplings

ex QED -  $m, e, Z_\psi, Z_\gamma$

$\lambda\phi \Rightarrow m, \lambda, Z_\phi$

All parameters in  $\mathcal{L}$  are dimensionless or positive mass dimension  $\times$

$\uparrow$   
 $e, \lambda$

$m^2 \phi^2, m \bar{\psi} \psi$

Field mass dimension  $[\phi] \sim m^1, [\partial\phi] \sim m^{3/2}, [A] = m^1, [\mathcal{L}] = m^4$

$\Rightarrow$  Field + derivative  $\underline{\underline{[O] \leq m^4}}$   $\leftarrow$

Why? expansion

$$\left[ 1 + \lambda \frac{1}{d-4} + \dots \right]$$

← don't generate  
higher dimensional  
operators

# Example

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \phi)^2 - m^2 \phi^2 - \frac{\lambda}{4} \phi^4 - \frac{g_6}{M^2} (\partial_\mu \phi)^2 \phi^2 - \frac{g_8}{M^4} (\partial_\mu \phi \partial^\mu \phi)^2$$

↑  
"renorm"
↑  
"non renormalizable"

One loop

$$X + \text{[Diagram: tadpole with loop]} = -6i\lambda + (6i\lambda)^2 \left[ \frac{1}{d-4} + \dots \right] = -6i\lambda_{ren}$$

$$+ \text{[Diagram: tadpole with loop and external legs]} + \frac{g_6^2}{M^4} f(m^2, m^2, p^2, p^4) \frac{1}{d-4}$$

ren λ
g<sub>6</sub>
g<sub>8</sub>
← can be renormalized

But need to renormalize  $g_8 \sim g_8^i \frac{1}{d-4}$ ,  $g_{10}$ ,  $g_{12} \dots$   
 $\Rightarrow$  more and more operators