

Introducing the Fields 2

Note Title

2/2/2010

Propagator

$$i D_F(x-x') = \langle 0 | T(\phi(x,t) \phi(x',t')) | 0 \rangle$$

$$= \int \frac{d^3 p}{(2\pi)^3} \frac{1}{2\omega} \left[\theta(t-t') e^{-ip \cdot (x-x')} + \theta(t'-t) e^{-ip \cdot (x'-x)} \right]$$

$\omega = \sqrt{p^2 + m^2}$

$$= \int \frac{d^4 q}{(2\pi)^4} i \frac{e^{-iq \cdot (x-x')}}{q^2 - m^2 + i\epsilon}$$

$\epsilon = 0^+$

$$\uparrow \quad \quad \quad \nwarrow$$

$d^3 q_0, d^3 q$ $q_0^2 - \vec{q}^2$

$$\int_{-\infty}^{\infty} dq_0 \frac{e^{-iq_0(t-t')}}{q_0^2 - \vec{q}^2 - m^2 + i\epsilon} = \oint_{\text{LHP (UHP)}} dq_0 \frac{e^{-iq_0(t-t')}}{q_0^2 - \dots} = -2\pi i \text{Res}(q_0 = \omega_q - i\epsilon)$$

$i\epsilon$ Physics

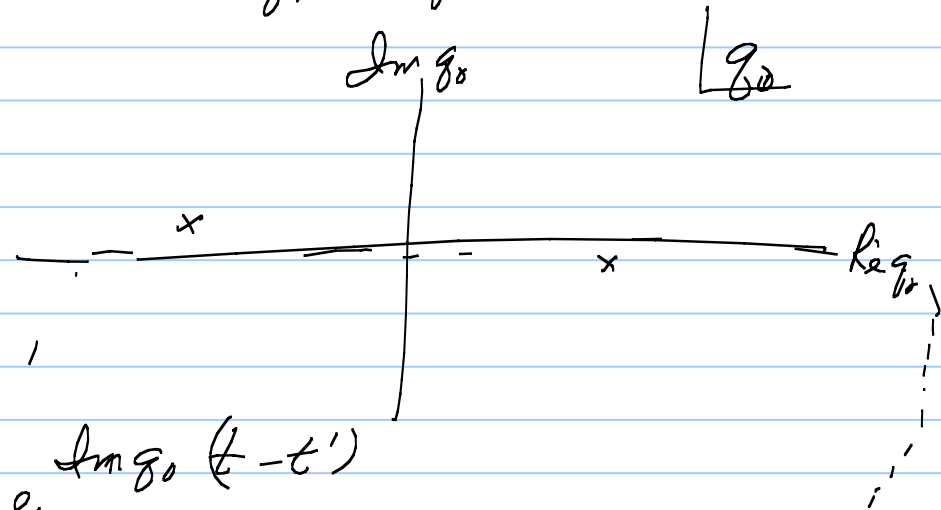
Poles $q_0^2 - \vec{q}^2 - m^2 + i\epsilon = 0$

$$q_0 = \sqrt{\vec{q}^2 + m^2} - i\epsilon$$

$$q_0 = -\sqrt{\vec{q}^2 + m^2} + i\epsilon$$

$$W_q = \sqrt{\vec{q}^2 + m^2}$$

$$q_0^2 = (\vec{q}^2 + m^2) - i\epsilon$$



To do contour integral

$$e^{-i q_0 (t-t')} = e^{-i \text{Re } q_0 (t-t')} e^{\text{Im } q_0 (t-t')}$$

if $t > t'$, $t-t' > 0$

$\text{Im } q_0 < 0$ on circle at $\infty \Rightarrow e^{-\infty} \Rightarrow$ LHP

if $t' > t$, $t-t' < 0$

$\text{Im } q_0 > 0 \Rightarrow$ UHP

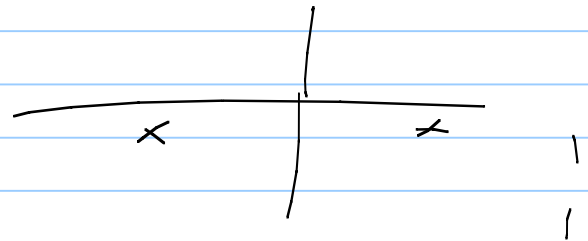
$$\frac{1}{q^2 - m^2 + i\epsilon} = \frac{1}{q_0 - W_q + i\epsilon} \frac{1}{q_0 + W_q - i\epsilon}$$

$$\begin{aligned}
 i D_F &= -2\pi i \int \frac{d^3q}{(2\pi)^3} \left[\theta(t-t') \frac{1}{2\omega_q} e^{-iE(t-t')} e^{i\vec{q}\cdot(\vec{x}-\vec{x}')} - \theta(t'-t) \frac{1}{-2\omega_q} e^{-i(-\omega_q)(t-t')} e^{i\vec{q}\cdot(\vec{x}-\vec{x}')} \right] \\
 &= \int \frac{d^3q}{(2\pi)^3} \frac{1}{2\omega_q} \left[\theta(t-t') e^{-i\omega_q(t-t')} e^{i\vec{q}\cdot(\vec{x}-\vec{x}')} + \theta(t'-t) e^{+i\omega_q(t-t')} e^{i\vec{q}\cdot(\vec{x}-\vec{x}')} \right] \checkmark
 \end{aligned}$$

Recall E+M Green functions for radiation (retarded G.F.)

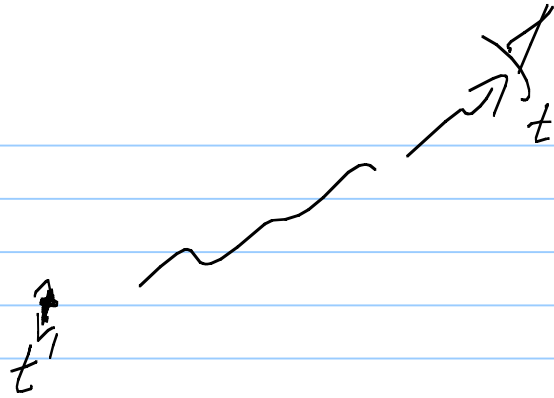
$$\frac{e^{-i\vec{q}\cdot(\vec{x}-\vec{x}')}}{(q_0 + i\epsilon) - \vec{q}^2}$$

$t > t'$ only



Physics

E+M



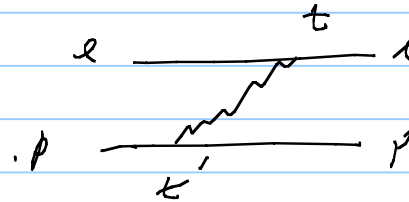
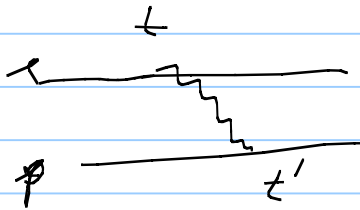
Want $t > t'$ only
for causality

QFT Interactions



← Feynman diagrams (no time implied)

Physics is



(time ordered pictures)

$$\text{Pert Theory} \quad \sum_I \frac{\langle f | V | I \rangle \langle I | V | i \rangle}{E - E_I}$$

↗ 2 intermediate states

Zero point energy

$$\sum_s \frac{1}{2} \hbar \omega_s (a(s) a^\dagger(s) + a^\dagger(s) a(s)) = \sum_s \hbar \omega_s a^\dagger(s) a(s) + \bar{E}_0 \quad \checkmark$$

$$E_0 = \sum_s \frac{1}{2} \hbar \omega_s = V \int \frac{d^3 p}{(2\pi\hbar)^3} \frac{1}{2} \hbar \omega_p \rightarrow \infty$$

$\delta^3(0)$

Various possibilities

1) really not infinite - "only N springs" - phonons, strings ...

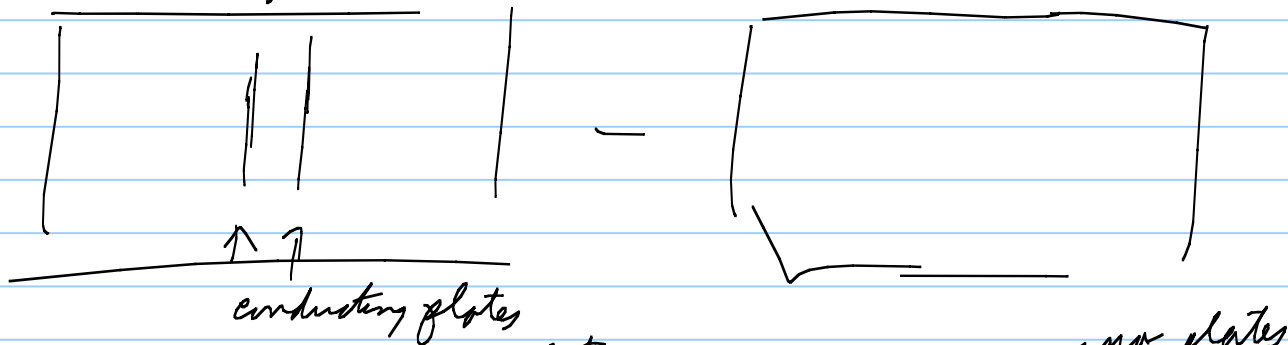
2) cancellation - bosons $E_0 > 0$ } "Supersymmetry" $E_0 \Rightarrow 0$
fermion $E_0 < 0$

3) gravity finite?

4) ?

Once finite \Rightarrow cosmological constant $\Lambda =$ energy of empty space
 if CMB are $\Lambda \Rightarrow \Lambda = (10^{-3} \text{ eV})^4$ density

Casimir effect



$$D \bar{E}_0 = \sum \frac{1}{2} \hbar \omega_i \text{ plates} - \sum \frac{1}{2} \hbar \omega_{no} \text{ no plates} \neq 0$$

$$\frac{F}{A} = - \frac{\pi^2 \hbar c}{240 d^4}$$

Complex scalar field $\phi, \phi^* \Rightarrow 2 \text{ D.O.F.}$

$$\mathcal{L} = \underline{\partial_\mu \phi^*} \partial^\mu \phi - V(\phi^* \phi)$$

$$V(\phi^* \phi) = \underline{m^2 \phi^* \phi} + \underline{\lambda (\phi^* \phi)^2}$$

Comments

1) ϕ, ϕ^* equiv to 2 real fields

$$\phi = \frac{1}{\sqrt{2}} (\phi_1 + i \phi_2)$$

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \phi_1 \partial^\mu \phi_1 + \partial_\mu \phi_2 \partial^\mu \phi_2) - \frac{m^2}{2} (\phi_1^2 + \phi_2^2) - \frac{\lambda}{4} (\phi_1^2 + \phi_2^2)^2$$

2) Symmetry $\phi \rightarrow e^{-i\theta} \phi$
 $\phi^* \rightarrow e^{+i\theta} \phi^*$ } phase "U(1)"

3) Eq of motion treat ϕ + ϕ^* as independent

$$\underline{\underline{\partial_\mu \frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^*}}} - \frac{\partial \mathcal{L}}{\partial \phi^*} = 0$$

$$\frac{\partial \mathcal{L}}{\partial \partial_\mu \phi^*} = \partial_\mu \phi \quad , \quad \frac{\partial \mathcal{L}}{\partial \phi^*} = -m^2 \phi - 2\lambda(\phi^* \phi) \phi$$

$$\square \phi - (-m^2 \phi - 2\lambda(\phi^* \phi) \phi) = 0$$

$$\underbrace{(\square + m^2)}_{\star} \phi = \underbrace{-2\lambda(\phi^* \phi) \phi}_{\text{interaction}}$$

Conserved current

$$J^m = (\rho, \vec{J}) \Rightarrow \partial_\mu J^m = 0 \equiv \frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot \vec{J}$$

$$J^m = i \phi^* \partial^m \phi - i (\partial^m \phi^*) \phi \equiv i \phi^* \overleftrightarrow{\partial}_m \phi$$

Check

$$\partial_\mu J^m = i \left[\underbrace{\partial_\mu \phi^* / \partial^m \phi}_{m^2 \phi^* \phi} + \underbrace{\phi^* \square \phi}_{m^2 \phi^* \phi} - (\square \phi^*) \phi - \partial_\mu \phi^* / \partial^m \phi \right] = 0$$

Like $E \times M$ conserved charge

$$Q = \int d^3x \underbrace{J^0(x, t)}_{\rho}$$

$$\frac{dQ}{dt} = \int d^3x \frac{\partial J^0}{\partial t} = \int d^3x (-\vec{\nabla} \cdot \vec{J}) = - \oint d\vec{A} \cdot \vec{J} \rightarrow 0 \quad \leftarrow \text{surface at } \infty$$

Calculate Q

$$Q = \int \frac{d^3p}{(2\pi)^3} (a^\dagger(p) a(p) - b^\dagger(p) b(p))$$

Interpretation - 2 sets of states

$$|\phi(p)\rangle = a^\dagger(p) |0\rangle$$

$$|\bar{\phi}(p)\rangle = b^\dagger(p) |0\rangle$$

← antiparticles

$$H |\phi(p)\rangle = \omega_p |\phi(p)\rangle, \quad H |\bar{\phi}(p)\rangle = \omega_p |\bar{\phi}(p)\rangle$$

$$Q |\phi(p)\rangle = + |\phi(p)\rangle, \quad Q |\bar{\phi}(p)\rangle = - |\bar{\phi}(p)\rangle$$