

Interactions 2

Note Title

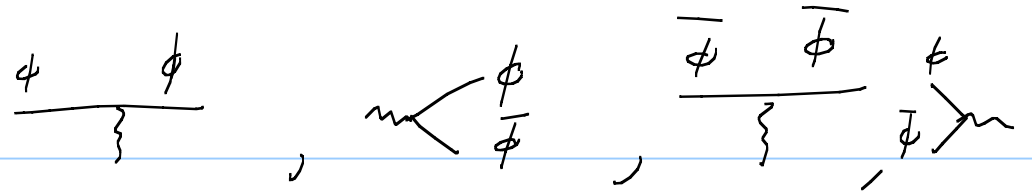
2/18/2010

$$i D_{\beta}(x-x') = \langle \beta | T(\phi(x) \phi(x')) | \beta \rangle$$

↑

| | | |
|-----------|-------------------------------|------------|
| 1st order | Order parameter discontinuous | ← |
| 2nd order | " " | continuous |

Crossing



$$\langle \phi(p_2) | J_\mu | \phi(p_1) \rangle = \dots g(p_1 + p_2)^\mu e^{-i(p_1 - p_2) \cdot x}$$

$$\langle \phi(p_2) \bar{\phi}(p_1) | J_\mu(x) | 0 \rangle = \dots g(p_2 - p_1)^\mu e^{+i(p_1 + p_2) \cdot x}$$

$$\langle \bar{\phi}(p_1) | J_\mu | \bar{\phi}(p_2) \rangle = \dots -g(p_1 + p_2)^\mu e^{-i(p_2 - p_1) \cdot x}$$

$$\langle 0 | J_\mu | \phi(p_1) \bar{\phi}(p_2) \rangle = \dots g(p_1 - p_2)^\mu e^{-i(p_1 + p_2) \cdot x}$$

Crossing rules

initial \Leftrightarrow final state

$p_i \rightarrow -p_i$

particle \leftrightarrow antiparticle

photon $E_\mu \leftrightarrow E^*$

Dirac $u \leftrightarrow v$

Dirac

$$\psi(x) = \int \frac{d^3 p}{(2\pi)^3} \left[e^{-ip \cdot x} u(p, s) b(p, s) + e^{+ip \cdot x} v(p, s) d^\dagger(p, s) \right]$$

$$J_\mu = g \bar{\psi} \gamma_\mu \psi$$

$$\langle f(p_2) | J_\mu | f(p_1) \rangle = g$$

$$= g \langle 0 | b(p_2) \underbrace{\int \frac{d^3 p'}{(2\pi)^3} \left[e^{+ip' \cdot x} \bar{u}(p') b^\dagger(p') \right]}_{(2\pi)^3 \delta^3(p' - p_2)} \gamma_\mu \underbrace{\int \frac{d^3 p}{(2\pi)^3} \left[e^{-ip \cdot x} u(p, s) b(p) \right]}_{(2\pi)^3 \delta^3(p - p_1)} b^\dagger(p_1, s) | 0 \rangle$$

$$= g \bar{u}(p_2) \gamma_\mu u(p_1) e^{-i(p_1 - p_2) \cdot x}$$

$$\begin{aligned}
 & \langle \bar{f}(p_2) | J_m | \bar{f}(p_1) \rangle \\
 &= g \langle 0 | d(p_2) \int \frac{d^3 p'}{(2\pi)^3} [e^{-i p' \cdot x} \bar{v}(p')] d(p') \gamma_m \int \frac{d^3 p}{(2\pi)^3} [\dots + e^{+i p \cdot x} d(p) v(p)] d^\dagger(p_1) | 0 \rangle
 \end{aligned}$$

↙ p₂

$$= -g \bar{v}(p_1) \gamma^m v(p_2) e^{-i(p_1 - p_2) \cdot x}$$

↑ ↙ ↘
 ! !
 antiparticle
 has opposite charge

At low velocity $\bar{u} \gamma_m u = g(1, \frac{\vec{p}}{m}) \underbrace{\chi^\dagger \chi}_1 = \bar{u} \gamma_m u$
 $\vec{p}_2 = \vec{p}_1$

Schewe

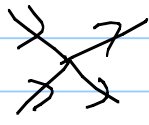
$$\langle \underline{f}(p_2) \bar{f}(p_1) | J_m | 0 \rangle = e^{+i(p_1 + p_2) \cdot x} g \bar{u}(p_2) \gamma_m u(p_1)$$

Other interaction

$$1) \lambda \phi^4 \Rightarrow \mathcal{L}_I = -\frac{\lambda}{4} \phi^4$$

like a (density)² interaction

$$\phi^2 \sim \psi^* \psi \Rightarrow \phi^4 \sim (\psi^* \psi) (\psi^* \psi)$$

in scattering 

$$2) \text{Complex scalar } \chi, \text{ real scalar } \phi \quad \mathcal{L}_I = -g \chi^* \chi \phi$$

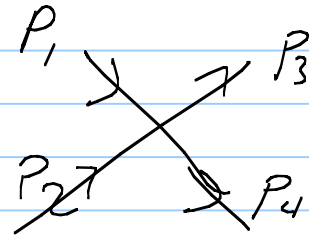
- like QED $\uparrow \sim$ charged particle $\uparrow \sim A_\mu$ $\frac{\chi}{\phi}$, $\frac{\bar{\chi}}{\phi}$

Samples

$$\mathcal{M} = \langle p_3, p_4 \mid \frac{\lambda}{4} \phi^4 \mid p_1, p_2 \rangle$$

$$= \frac{\lambda}{4} \frac{e^{-i(p_1+p_2-p_3-p_4) \cdot X}}{\sqrt{2w_1 2w_2 2w_3 2w_4}} \leftarrow P$$

↑ counting factor



To get P

4 way to annihil. p_1

3 way " p_2

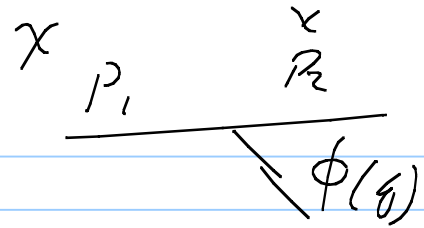
2 way to create p_3

1 way " p_4

$$P = 4!$$

$$\mathcal{M} = \left[\frac{e^{-i(\quad) \cdot X}}{\sqrt{2w_1 \dots 2w_4}} \right] 6\lambda$$

$$\begin{aligned}
 2) \quad \mathcal{M} &= \langle \chi(p_2) \phi(v) | g \chi^\dagger \chi \phi | \chi(p_1) \rangle \\
 &= \left[\frac{e}{\sqrt{2}} \right] g
 \end{aligned}$$



Perturbation theory

- Plan
- 1) Time Development operator
 - 2) Calculate transition amplitudes
 - 3) Feynman rules

Recall Interaction Picture $H = H_0 + H_I$

- 1) Basis states due to H_0
- for no particle states $|P_i\rangle$
- 2) Time independent if $H_I = 0$
- 3) States change due to interaction

$$|\psi(t)\rangle = U_I(t, t_0) |\psi(t_0)\rangle$$

↑ Time Development Op

$$U_I(t, t_0) = T e^{-i \int_{t_0}^t dt' H_I(t')}$$

Conceptual path

$$i \frac{d}{dt} |\psi_I(t)\rangle = H_I(t) |\psi_I(t)\rangle$$

Integral eq $|\psi_I(t)\rangle - |\psi_I(t_0)\rangle = -i \int dt' H_I(t') |\psi_I(t')\rangle$

Iterate $|\psi_I(t)\rangle = |\psi_I(t_0)\rangle - i \int_{t_0}^t dt' H_I(t') \left[|\psi_I(t_0)\rangle - i \int_{t_0}^{t'} dt'' H_I(t'') |\psi_I(t'')\rangle \right]$

$$= \left[1 - i \int_{t_0}^t dt' H_I(t') + (-i)^2 \int_{t_0}^t dt' \int_{t_0}^{t'} dt'' H_I(t') H_I(t'') + \dots \right] |\psi_I(t_0)\rangle$$

Change of variables $\int_{t_0}^t dt' \int_{t_0}^{t'} dt'' H_I(t') H_I(t'') = \int_{t_0}^t \int_{t_0}^t dt' dt'' T(H_I(t') H_I(t''))$

Compact notation $[\dots] = U_I(t, t_0) = T e^{-i \int_{t_0}^t dt' H_I(t')}$

Example $\mathcal{L}_I = -\frac{\lambda}{4} \phi^4$

$t \rightarrow -\infty \quad |p_1, p_2\rangle = |\psi_I^i\rangle$

Time develop $t \rightarrow +\infty \Rightarrow U_I(\infty, -\infty) |p_1, p_2\rangle$

Amplitude for finding $|\psi_f\rangle = |p_3, p_4\rangle$

$T_{fi} = \langle p_3, p_4 | U_I(\infty, -\infty) | p_1, p_2 \rangle$

\uparrow S matrix $= T \exp \int_{-\infty}^{+\infty} dt H_I(t) = T \exp \left\{ -i \int dt d^3x \mathcal{H}_I(x, t) \right\}$
 $= T \exp \left\{ +i \int d^4x \mathcal{L}_I(x, t) \right\}$

our examples $\mathcal{H}_I = -\mathcal{L}_I$

To first order in λ

$$T_{fi} = \langle p_3, p_4 | \left[1 - i \frac{\lambda}{4} \int d^4x \phi^4(x,t) + \dots \right] | p_1, p_2 \rangle$$

$$= -i \int d^4x \left(\frac{\lambda}{4} \times 4! \right) \frac{e^{-i(p_1 + p_2 - p_3 - p_4) \cdot x}}{\sqrt{2\omega_1} \sqrt{2\omega_2} \sqrt{2\omega_3} \sqrt{2\omega_4}}$$

$$= -i \frac{\delta^4(p_1 + p_2 - p_3 - p_4) \mathcal{M}}{\sqrt{2\omega_1} \dots \sqrt{2\omega_2}}$$

$$\mathcal{M} = 6\lambda \quad \Rightarrow \quad -i\mathcal{M} = -6i\lambda$$