

Basics I

1/19/10

Note Title

1/19/2010

1905 - Quantum

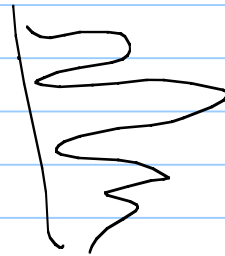
- field \rightarrow particles
- annihilated

Particles \rightarrow fields

Created

$e^+ e^- \rightarrow \mu^+ \mu^-$

QFT completion of QM



Course

- basics

- for all areas

- rel. notations

Books -

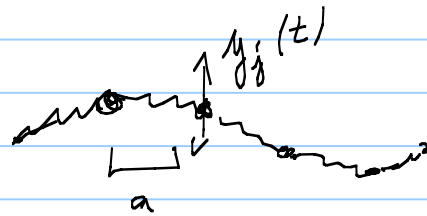
HW.

Recording

Office hours - all, Google calendar

Construct a field — $\phi(x, t)$

$y_j(t)$



$$L = \sum_j \left[\frac{1}{2} m \dot{y}_j^2 - \frac{1}{2} k (y_{j+1} - y_j)^2 \right] = L(y_j, \dot{y}_j)$$

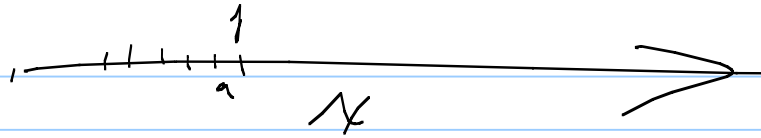
$$S = \int dt L$$

Classically

$$p_j = \frac{\partial L}{\partial \dot{y}_j} = m \dot{y}_j$$

$$H = \sum_j p_j \dot{y}_j - L = H(p_j, y_j) = \sum_j \left[\frac{p_j^2}{2m} + \frac{1}{2} k (y_{j+1} - y_j)^2 \right]$$

Continuum



$$N = a j$$

$$\sum_j a = \int dx$$

$$y_j(t) = \frac{1}{\sqrt{ak}} \phi(x, t)$$

$$\sum_j \frac{1}{2} k a^2 \left(\frac{y_{j+1} - y_j}{a} \right)^2 = \sum_j a \frac{1}{a} \frac{1}{2} k a^2 \frac{1}{ka} \left(\frac{\partial \phi}{\partial x} \right)^2 = \int dx \frac{1}{2} \left(\frac{\partial \phi}{\partial x} \right)^2$$

$$\sum_j \frac{1}{2} m \dot{y}_j^2 = \int dx \frac{1}{2} \frac{1}{v^2} \left(\frac{\partial \phi}{\partial t} \right)^2, \quad v = \sqrt{\frac{ka^2}{m}}$$

$$S = \int dt dx \left[\frac{1}{2} \frac{1}{v^2} \left(\frac{\partial \phi}{\partial t} \right)^2 - \frac{1}{2} \left(\frac{\partial \phi}{\partial x} \right)^2 \right] = \int dt dx \mathcal{L}(\phi, \dot{\phi}, \frac{\partial \phi}{\partial x})$$

\mathcal{L} density

$$L = \int dx \mathcal{L}$$

Euler Lagrange eq.

$$\phi(x,t) = \bar{\phi}(x,t) + \delta\phi(x,t)$$

↑ classical solution

↙ $\delta\phi(x,t) = 0$ on any boundaries
at infinity

$$\delta S = 0 = \int dt dx \left[\frac{1}{v^2} \frac{\partial \bar{\phi}}{\partial t} \frac{\partial \delta\phi}{\partial t} - \frac{\partial \bar{\phi}}{\partial x} \frac{\partial \delta\phi}{\partial x} \right]$$

↑ PLA

$$= \int dt dx \left[-\frac{1}{v^2} \frac{\partial^2 \bar{\phi}}{\partial t^2} + \frac{\partial^2 \bar{\phi}}{\partial x^2} \right] \delta\phi(x,t)$$

↑ int by parts

for any $\delta\phi(x,t)$

$$\left[\frac{1}{v^2} \frac{\partial^2}{\partial t^2} - \frac{\partial^2}{\partial x^2} \right] \bar{\phi}(x,t) = 0$$

wave eq.

Hamiltonians

$$p_j = \frac{\partial L}{\partial \dot{y}_j(t)} \quad \Rightarrow \quad \frac{\partial \mathcal{L}}{\partial \phi(x,t)} \equiv \pi(x,t) = \frac{1}{v^2} \frac{\partial \phi}{\partial t}$$

$$H = \sum_i \left[\frac{p_i^2}{2m} + \frac{1}{2} k (y_{i+1} - y_i)^2 \right] = \int dx \left[\underbrace{\frac{1}{2} v^2 \left(\frac{\partial \phi}{\partial t} \right)^2}_{\frac{1}{2} \pi^2 v^2} + \frac{1}{2} \left(\frac{\partial \phi}{\partial x} \right)^2 \right] \mathcal{H}$$

Field theory

$$H = \int dx \mathcal{H}$$

$$\mathcal{H} \equiv \pi \dot{\phi} - \mathcal{L} = \left[\frac{1}{2} \pi^2 v^2 + \frac{1}{2} \left(\frac{\partial \phi}{\partial x} \right)^2 \right] = \mathcal{H}(\pi, \phi)$$

Normalization

$$P_i = m \dot{y}_i = \frac{m}{\sqrt{\hbar a}} \dot{\phi}(x, t) = \frac{1}{\nu^2} a \sqrt{\hbar a} \dot{\phi} = a \sqrt{\hbar a} \pi(x, t)$$

$$y_i(t) = \frac{1}{\sqrt{\hbar a}} \phi(x, t)$$



Quantization

$$[y_i, p_{j'}] = i\hbar \delta_{i,j'}$$

$$= \left[\frac{1}{\sqrt{\hbar a}} \phi(x,t), \sqrt{\hbar a} \pi(x',t) \right]$$

$$[\phi(x,t), \pi(x',t)] = i\hbar \frac{\delta_{j,j'}}{a} = i\hbar \delta(x-x')$$

↖ Equal time commutation rule

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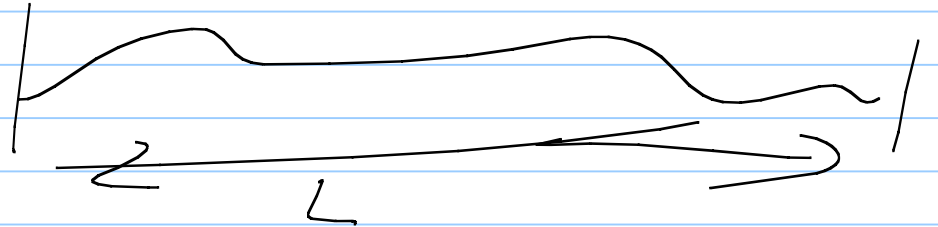
Correspondence

$$\sum_{j'} \delta_{j,j'} = 1 = \left(\sum_{j'} a \right) \frac{\delta_{j,j'}}{a} = \int dx' \delta(x-x') = 1$$

$\underbrace{\qquad}_{\delta(x-x')} = \lim_{a \rightarrow 0} \frac{\delta_{j,j'}}{a} \quad x = ja$

Quanta - discrete notations

- Box normalization
- BC not important
- energy levels discrete



Solutions

$$\left[\frac{1}{\nu^2} \frac{\partial^2}{\partial t^2} - \frac{\partial^2}{\partial x^2} \right] \phi(x, t) = 0 \quad \phi = N \left[e^{-i(\frac{\omega_s}{\nu} t - k_s x)} + c.c. \right]$$

$$\frac{\omega_s^2}{\nu^2} - k_s^2 = 0 \quad \Rightarrow \quad \omega_s = \nu k_s$$

General solution

$$\phi(x, t) = \sum_s N_s \left(e^{-i(\omega_s t - k_s x)} a(s) + e^{i(\omega_s t - k_s x)} a^\dagger(s) \right)$$

Conditions

$$[a(s), a(s')] = 0$$

$$[a(s), a^+(s')] = \delta_{ss'}$$

$$N_s = \sqrt{\frac{\hbar N^2}{2\omega_s L}}$$

$$\Pi(x, t) = \frac{1}{N^2} \frac{\partial \phi}{\partial t} = \sum_{s'} N_{s'} \frac{-i\omega_{s'}}{N^2} \left(e^{-i(\omega_{s'}t - k_{s'}x)} a(s') - e^{+i(\omega_{s'}t - k_{s'}x)} a^+(s') \right)$$

Check:

$$\begin{aligned} [\phi(x, t), \Pi(x', t)] &= \sum_{s, s'} N_s N_{s'} \left(\frac{-i\omega_{s'}}{N^2} \right) \left[\left(e^{-i(\omega_s t - k_s x)} a(s) + e^{+i(\omega_s t - k_s x)} a^+(s) \right), \left(e^{-i(\omega_{s'} t - k_{s'} x')} a(s') - e^{+i(\omega_{s'} t - k_{s'} x')} a^+(s') \right) \right] \\ &= \sum_s N_s^2 \underbrace{\left(\frac{+\omega_s}{N^2} \right)}_{\frac{i\hbar}{L}} 2 e^{i k_s (x - x')} = i\hbar \sum_s \frac{1}{L} e^{i k_s (x - x')} \\ &= i\hbar \delta(x - x') \quad \checkmark \end{aligned}$$

Hamiltonian

$$H = \int dx \frac{1}{2} \left[\frac{1}{\mu^2} \left(\frac{\partial \phi}{\partial t} \right)^2 + \left(\frac{\partial \phi}{\partial x} \right)^2 \right]$$

⋮

$$= \sum_s \hbar \omega_s \left(a^\dagger(s) a(s) + \frac{1}{2} \right)$$

$\underbrace{\hspace{2cm}}_{\# \text{ operators}} \quad \leftarrow E_0$

$$E = n \hbar \omega_s \quad \leftarrow \underline{\text{quanta}}$$

